#### **General Disclaimer**

#### One or more of the Following Statements may affect this Document

- This document has been reproduced from the best copy furnished by the organizational source. It is being released in the interest of making available as much information as possible.
- This document may contain data, which exceeds the sheet parameters. It was furnished in this condition by the organizational source and is the best copy available.
- This document may contain tone-on-tone or color graphs, charts and/or pictures, which have been reproduced in black and white.
- This document is paginated as submitted by the original source.
- Portions of this document are not fully legible due to the historical nature of some
  of the material. However, it is the best reproduction available from the original
  submission.

Produced by the NASA Center for Aerospace Information (CASI)

### Technical Memorandum 33-785

## Disturbing Effects of Attitude Control Maneuvers on the Orbital Motion of the Helios Spacecraft

(NASA-CR-148840) DISTURBING EFFECTS OF
ATTITUDE CONTRC MANEUVERS ON THE OFBITAL
MOTION OF THE F LICS SPACECRAFT (Jet
Propulsion Lab.) 44 p HC \$4.00 CSCL 22A
G3/13



N76-33244

Unclas

07966



JET PROPULSION LABORATORY

CALIFORNIA INSTITUTE OF TECHNOLOGY

PASADENA, CALIFORNIA

September 15, 1976

$\overline{\cdot}$	22.705	2 Comment	Accession No.	3. Recipient's Catalog No.
١.	. Report No. 33-785  2. Government Accession No.			
4. Title and Subtitle			5. Report Date September 15, 19	
	ON THE ORBITAL MOTION OF THE HELIOS SPACECRAFT  7. Author(s)  R. M. Georgevic			6. Performing Organization Code
7.			8. Performing Organization Report No.	
9. Performing Organization Name and Address			10. Work Unit No.	
	JET PROPULSION LABORATORY California Institute of Technolog		gy	11. Contract or Grant No. NAS 7-100
4800 Oak Grove Drive Pasadena, California 91103  12. Sponsoring Agency Name and Address  NATIONAL AERONAUTICS AND SPACE ADMINISTR Washington, D.C. 20546			13. Type of Report and Period Cover	
		Technical Memorandum		
		RATION	14. Sponsoring Agency Code	
	and updated by a series of unbalanced jet force spacecraft's center of motion, its magnitude a purposes of the orbit of shows how the component fixed reference frame	s of attitude ces which produce mass. This reand direction. determination of the disturbance be easily con be easily con the can be easily con the disturbance of the	ontrol maneue an addition port examine In addition f the spaced rbing acceler	deceraft has been maintained evers, by means of a sequence onal disturbed motion of the esthe character of this ato this, for practical eraft, a computer program eration in the spacecraft—se program is given as an
	appendix to this repor-	t.		
17.	Key Words (Selected by Author(s	3))	18. Distributio	on Statement
	Spacecraft Communication and Tracking	514.0	Unc]	Lassified Unlimited

20. Security Classif. (of this page)

Unclassified

21. No. of Pages

38

22. Price

19. Security Classif. (of this report)

Unclassified

### Technical Memorandum 33-785

## Disturbing Effects of Attitude Control Maneuvers on the Orbital Motion of the Helios Spacecraft

R. M. Georgevic

JET\_PROPULSION LABORATORY
CALIFORNIA INSTITUTE OF TECHNOLOGY
PASADENA, CALIFORNIA

September 15, 1976

#### PREFACE

The work described in this report was performed by the Mission Analysis Division of the Jet Propulsion Laboratory.

#### CONTENTS

Introduction	1
The Three Reference Frames	4
The Position of the Attitude Control Valve	6
Types of Maneuvers	10
Components of the Force in the Equatorial, Space-Fixed Reference Frame	11
Integration of Equations of Motion	14
Geometry of Motion	15
The Motion During One Cycle (16 Pulses)	17
One Example	19
The Computer Program	21
The Implementation of the Perturbative Acceleration in the Orbit Determination Program	21
Accelerations in the Sun-Canopus Oriented System	24
Appendix Computer Program for Calculation of Helios Maneuvers	29
References	38
Bibliography	38

ING PAGE BLANK NOT FILM

TABLES		
1.	Helios A spacecraft maneuvers	3
2.	Types of maneuvers	13
3.	Accumulated disturbing effects in coordinates during maneuvers	22
4.	Helios attitude control maneuvers, disturbing accelerations .	27
FIGURES		
1.	Ecliptic reference systems	5
2	Spacecraft body-fixed pitch-yaw-roll system	7
3.	Geometry of the maneuvers	9
4.	Schematic view of the maneuvers	12

#### ABSTRACT

The position of the spin axis of the Helios A spacecraft has been maintained and updated by a series of attitude control maneuvers, by means of a sequence of unbalanced jet forces which produce an additional disturbed motion of the spacecraft's center of mass. This report examines the character of this motion, its magnitude and direction. In addition to this, for practical purposes of the orbit determination of the spacecraft, a computer program shows how the components of the disturbing acceleration in the spacecraft-fixed reference frame can be easily computed. The program is given as an appendix to this report.

# DISTURBING EFFECTS OF ATTITUDE CONTROL MANEUVERS ON THE ORBITAL MOTION OF THE HELIOS SPACECRAFT

#### R. M. Georgevic

#### In roduction:

The unbalanced attitude control nozzle firings on the spin-stabilized Helios spacecraft during the sequence of attitude control maneuvers are producing disturbing effects on the orbital motion of the spacecraft. This additional motion of the center of mass of the spacecraft occurs only when the jet force of the nozzle is not counteracted by another jet force of the same magnitude and opposite direction, thus producing a couple which causes a purely rotational motion of the spacecraft.

The jet force is not continuous: it starts when the Sun sensor activates the nozzle and lasts for one quarter of the rotational period of the space-craft, i.e. for one fourth of a second and then stops, to start again at the beginning of the next second. Each cycle has sixteen such pulses; between each two cycles the nozzle is inactive and produces no force. The maneuvers of the Helios A spacecraft are shown in Table 1.

In further text we shall entirely disregard the rigid-body motions of the spacecraft and examine only the effects of jet forces on its orbital motion, considering the spacecraft as a point-mass. Still, to simplify a rather complicated analysis of the motion under the influence of these forces

combined with the motion under other forces acting on the spacecraft, we shall introduce a few restrictions. They are as follows:

- 1. We shall observe the effects of jet forces during one cycle (16 pulses) separated from the effects of other forces. This can be done because of the short duration of each cycle (16 seconds) and the small magnitude of the jet force (1 Nt).
- The spinning period of the spacecraft will be assumed constant,
   e.g. the spin rate will always be 60 rpm.
- 3. The fundamental plane (xy-plane) of the rotating, spacecraft-fixed system of reference will be the ecliptic plane of 1950.0. This can be assumed because the position of the spin-axis of Helios is almost perfectly colinear with the normal to the ecliptic plane (or at least within tolerable limits).
- During one cycle of sixteen pulses the spacecraft will be considered stationary, e.g.

$$\overline{\mathbf{r}}(\mathbf{t}) = \overline{\mathbf{r}}(\mathbf{t}_0), \overline{\mathbf{v}}(\mathbf{t}) = 0, \mathbf{t}_0 \leq \mathbf{t} \leq \mathbf{t}_0 + 16 \text{ sec.},$$

where  $t_0$  is the beginning of the cycle.

5. It will also be assumed that, at the beginning of each cycle, the y-axis (pitch-axis of Helios carrying the Sun sensor) of the rotating spacecraft-fixed reference frame is pointing at the Sun.

Table 1. Helios A spacecraft maneuvers

DAY OF YEAR				,	
1975	HR	MIN	SEC	TYPE	DUKATION
41	07	29	00	RP	16 Pulses
	07	38	00	RP	"
	12	06	00	RN	"
	12	15	00	RN	"
	12	24	00	RN	"
	12	33	00	RN	"
	13	56	00	PN	
	14	24	00	PN PN	"
	14 15	52 20	00 00	PN	16 Pulses
	15	20	00	FN	10 ruises
63	07	55	00	RN	16 Pulses
0.5	C8	09	21	RN	"
	08	22	22	RN	"
	08	35	43	RN	"
	_9	36	59	PN	"
	10	06	44	PN	"
	14	12	45	PN	"
	14	26	45	SPUP	4 Revolutions
	15	17	39	PN	16 Pulses
	15	49	51	PN	"
	16	15	53	PN	16 Pulses
80	07	29	00	RN	16 Pulses
00	07	50	00	RN	"
	08	30	00	RN	"
	08	50	00	RN	"
	09	10	00	RN	"
	09	30	00	RN	u u
	09	50	00	RN	16 Pulses
					16 8 1
102	09	09	13	PN	16 Pulses
	09	40	00	PN	,,
	09	50	00	PN	"
	10	00	00	PN	"
	10	10	00 00	PN PN	"
	10	20 30	00	PN	"
	10	50 50	00	RN	,,
	10 11	00	00	RN	"
	11	10	00	RN	"
	11	20	00	RN	16 Pulses
	11				1

RP = ROLL POSITIVE

RN = ROLL NEGATIVE

PN = PITCH NEGATIVE

#### The Three Reterence Frames

Without specifying the positions of their origins, we shall introduce two ecliptic reference systems: the first one, XYZ is ecliptic, spacefixed system, with its X-axis toward the vernal equinox point; the second is the system rotating with the spacecraft at 60 rpm. Both systems are shown on Figure 1. The relationship between the two systems is given by

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} \cos\phi & -\sin\phi & 0 \\ \sin\phi & \cos\phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} , \qquad (1)$$

where  $\phi$  is the angle shown on Figure 1. It is obvious that, due to the spacecraft's rotation, the angle  $\phi$  is proportional to  $\omega t$ .

Introducing now the space-fixed equatorial frame of reference,  $\mathbf{X}_{Q}$  ,  $\mathbf{Y}_{Q}$  ,  $\mathbf{Z}_{Q}$  , we find that

$$\begin{pmatrix} x_{Q} \\ Y_{Q} \\ Z_{Q} \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \varepsilon & \sin \varepsilon \\ 0 & \sin \varepsilon & \cos \varepsilon \end{pmatrix} \begin{pmatrix} x \\ Y \\ Z \end{pmatrix} , \qquad (2)$$

and, combining (1) and (2), we obtain

$$\begin{pmatrix} \mathbf{X}_{\mathbf{Q}} \\ \mathbf{Y}_{\mathbf{Q}} \\ \mathbf{Z}_{\mathbf{Q}} \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \varepsilon & -\sin \varepsilon \\ 0 & \sin \varepsilon & \cos \varepsilon \end{pmatrix} \begin{pmatrix} \cos \varphi & -\sin \varphi & 0 \\ \sin \varphi & \cos \varphi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \mathbf{x} \\ \mathbf{y} \\ \mathbf{z} \end{pmatrix} ,$$

or

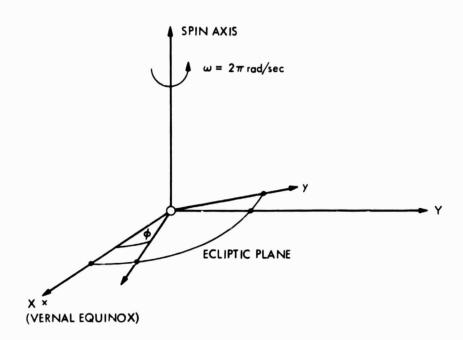


Fig. 1. Ecliptic reference systems

$$\begin{pmatrix} x_{Q} \\ Y_{Q} \\ z_{Q} \end{pmatrix} = \begin{pmatrix} \cos\phi & -\sin\phi & 0 \\ \cos\varepsilon & \sin\phi & \cos\varepsilon & \cos\phi & -\sin\varepsilon \\ \sin\varepsilon & \sin\phi & \sin\varepsilon & \cos\phi & \cos\varepsilon \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} . \quad (3)$$

#### The Position of the Attitude Control Valve

The pitch and roll motions of Helios are determined with respect to another reference frame which, for the time of duration of one maneuvering cycle, can be considered stationary  $(x_1 \ y_1 \ z_1 - \text{system})$ . The  $x_1y_1$ -plane of this system coincides with the ecliptic plane and the  $y_1$ -axis (pitch axis) is always in the spacecraft-sun direction (Figure 2). The angle  $\lambda_s$  is the longitude of the Sun observed from the spacecraft, which is

$$\lambda_{s} = 180^{\circ} + \lambda,$$

where  $\lambda$  is the heliocentric longitude of the spacecraft. Hence,

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} -\sin \lambda & -\cos \lambda & 0 \\ \cos \lambda & -\sin \lambda & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x_1 \\ y_1 \\ z_1 \end{pmatrix} \tag{4}$$

Combining Equations (1) and (4) we find

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} \sin(\phi - \lambda) & -\cos(\phi - \lambda) & 0 \\ \cos(\phi - \lambda) & \sin(\phi - \lambda) & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x_1 \\ y_1 \\ z_1 \end{pmatrix}$$
(5)

Measuring time since the moment when the two systems coincide, we obtain

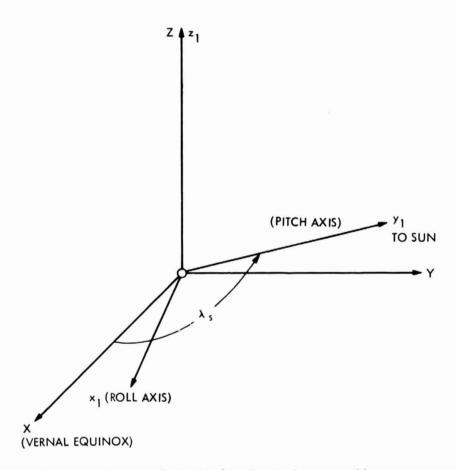


Fig. 2. Spacecraft body-fixed pitch-yaw-roll system

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} \cos \omega t & \sin \omega t & 0 \\ -\sin \omega t & \cos \omega t & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x_1 \\ y_1 \\ z_1 \end{pmatrix} . \tag{6}$$

Comparison of Equations (5) and (6) yields

$$\phi = \lambda + \omega t + \frac{\pi}{2} \qquad , \tag{7}$$

and, for t = 0  $(y \equiv y_1)$ ,

$$\phi_0 = \lambda + \frac{\pi}{2} .$$

The position of the precession valve is shown on Figure 3. The angle  $\theta$  is approximally  $38^{\circ}$ . The nozzle is located on the lowest brim of the truncated cone of the Helios body. The line drawn from the center of mass of the spacecraft (C.M. on Figure 3) to the nozzle is perpendicular to the axis of symmetry of the nozzle – the line of action of the jet force. The nozzle lies in a plane parallel to the yz-plane of the body-fixed rotating reference frame. The magnitude of the jet force is

$$F = |\dot{m}|v_e \approx 1 \text{ Nt}$$

where  $\dot{m}$  is the gas flow and  $v_e$  is the nozzle velocity of the expelled gas. The magnitude of the acceleration of this force is (mass of the spacecraft is m = 356.9 kg):

$$a = 2.802 \times 10^{-6} \text{ km/sec}^2$$
.

To estimate the influence of only one nozzle firing in the duration of one fourth of a second, on the motion of the center of mass of the spacecraft, we find first the cange in velocity after the firing, which is

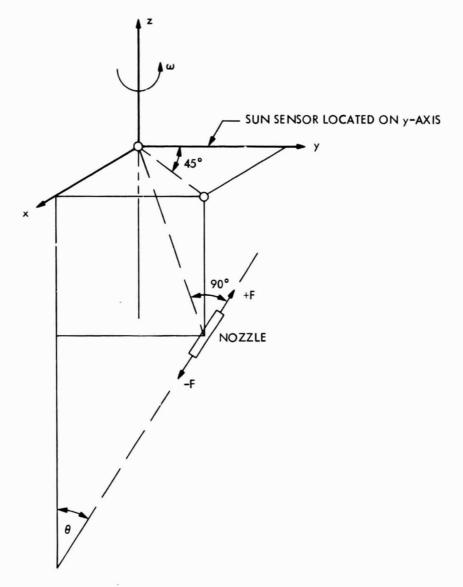


Fig 3. Geometry of the maneuvers

$$\Delta v = 7.0 \times 10^{-7} \text{ km/sec.}$$

Disregarding the orbital motion of the spacecraft, this change in velocity produces a change in the spacecraft's position in an amount of

$$\left|\Delta \bar{\mathbf{r}}\right| = 60.5 \text{ m/day.}$$

#### Types of Maneuvers

From Table 1. we see that there are four different types of maneuvers; they are:

RP - Positive roll rotation abouts the +x-axis

RN - Negative roll rotation about the +x-axis.

PN - Negative pitch rotation about the +y-axis

SPUP - (Spin up) Positive rotation about the +z-axis: increase of the spin rate.

The last type of maneuver does not produce any motion of the center of mass of the spacecraft. The rotation rate is changed by a couple of jet valves which operate in a plane perpendicular to the spin-axis and generate a couple; therefore, assuming that the jet forces on both exhaust valves are balanced (or close to being equal), there is no excess force acting on the center of mass.

To create a rotational motion purely about the  $x_1$ -axis (roll motion), the precession valve is activated when the angle between the y-axis and the  $y_1$ -axis (Sun direction) is  $45^\circ$  (actually -45°). The jet force creates also a pitch motion but, since the valve operates only for a quarter of a second (one quadrant), during the first one-eighth of a second the pitch motion is negative and during the next one-eighth of a second it is positive and

same in magnitude. Hence the total pitch motion over the complete firing period is zero and the outcome is a purely roll motion. The same is done to produce a pure rotation about the  $y_1$ -axis, by cancelling out the roll motion.

Since we are not concerned with rotational motions of the spacecraft, we need to know only in which direction the jet force is acting for each particular type of maneuver. From Figure 3 we can see that the PN and RP motions are produced by a +F force (out of ecliptic plane to the northern hemisphere), and that RN motion is produced by a -F force (out of the ecliptic plane, south).

RP, RN and PN motions are schematically shown on Figure 4, and listed in Table 2.

#### Components of the Force in the Equatorial, Space-Fixed Reference Frame

From Figure 3 we see that the components of the jet force, along the axes of the rotating, body-fixed reference frame are

$$\begin{pmatrix} 0 \\ f \sin \theta \\ f \cos \theta \end{pmatrix},$$

where

The components of the force in the ecliptic, space-fixed reference frame are

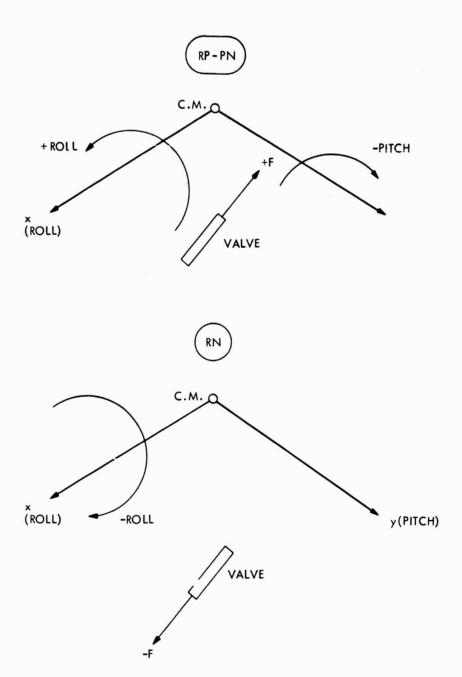


Fig. 4. Schematic view of the maneuvers

$$\begin{pmatrix} \mathbf{F}_1 \\ \mathbf{F}_2 \\ \mathbf{F}_3 \end{pmatrix} = \begin{pmatrix} \cos\phi & -\sin\phi & 0 \\ \sin\phi & \cos\phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 \\ f \sin\theta \\ f \cos\theta \end{pmatrix}$$

and in the space-fixed equatorial reference frame,

$$\begin{pmatrix} \mathbf{F}_{\mathbf{x}} \\ \mathbf{F}_{\mathbf{y}} \\ \mathbf{F}_{\mathbf{z}} \end{pmatrix} = \begin{pmatrix} \cos\varphi & -\sin\varphi & 0 \\ \cos\varepsilon\sin\varphi & \cos\varepsilon\cos\varphi & -\sin\varepsilon \\ \sin\varepsilon\sin\varphi & \sin\varepsilon\cos\varphi & \cos\varepsilon \end{pmatrix} \begin{pmatrix} 0 \\ \mathbf{f}\cos\theta \\ \mathbf{f}\cos\theta \end{pmatrix} ,$$

or

$$F_{x} = -f \sin\theta \sin\phi,$$

$$F_{y} = f(\cos\epsilon \sin\theta \cos\phi - \sin\epsilon \cos\theta),$$

$$F_{z} = f(\sin\epsilon \sin\theta \cos\phi + \cos\epsilon \cos\theta).$$
(8)

where the angle  $\phi$  is given explicitly in Equation (7).

Table 2. Types of maneuvers

Maneuver Type	Means	Force	
RP	Positive Roll Motion	+F	
RN	Negative Roll Motion	-F	
PN	Negative Pitch Motion	+F	

#### Integration of Equations of Motion

Since the angular orbital motion of the spacecraft, at the rate of  $\dot{\lambda}=3.8\times10^{-7}$  rad/sec is negligible in comparison with the rotation rate  $\omega=2\pi$  rad/sec, we shall consider  $\lambda$  as constant. Equations of motion are

$$\begin{pmatrix} \mathbf{X}_{\mathbf{Q}} \\ \mathbf{X}_{\mathbf{Q}} \\ \mathbf{Y}_{\mathbf{Q}} \\ \mathbf{Z}_{\mathbf{Q}} \end{pmatrix} = \frac{\mathbf{f}}{\mathbf{m}} \begin{pmatrix} -\sin\theta & \sin\phi \\ \cos\epsilon & \sin\theta & \cos\phi & -\sin\epsilon & \cos\theta \\ \sin\epsilon & \sin\theta & \cos\phi & +\cos\epsilon & \cos\theta \end{pmatrix} ,$$

where  $\phi = \lambda + \omega t + \frac{\pi}{2}$ . At the beginning of the first valve firing t = 0 and, at the end of the first pulse  $\omega t = \frac{\pi}{2}$ . Also,  $d\phi = \omega dt$ . The first integration yields  $(\phi_0 = \lambda + \frac{\pi}{2})$  the change of the velocity of the center of mass of the spacecraft due to the action of the jet force in the form

$$\begin{pmatrix} \Delta \dot{\mathbf{x}}_{\mathrm{Q}} \\ \Delta \dot{\mathbf{y}}_{\mathrm{Q}} \end{pmatrix} = \frac{\mathrm{f}}{\mathrm{m}\omega} \begin{pmatrix} -(\cos\phi - \cos\phi_{\mathrm{o}})\sin\theta \\ \cos\epsilon(\sin\phi - \sin\phi_{\mathrm{o}})\sin\theta - (\phi - \phi_{\mathrm{o}})\sin\epsilon\cos\theta \\ \sin\epsilon(\sin\phi - \sin\phi_{\mathrm{o}})\sin\theta + (\phi - \phi_{\mathrm{o}})\cos\epsilon\cos\theta \end{pmatrix},$$

where  $\phi_0 \le \phi \le \frac{\pi}{2}$ , and  $\phi_0 = \lambda + \frac{\pi}{2}$ . Finally, we find

$$\begin{pmatrix} \Delta \dot{\mathbf{x}}_{\mathbf{Q}} \\ \Delta \dot{\mathbf{Y}}_{\mathbf{Q}} \end{pmatrix} = \frac{\mathbf{f}}{m\omega} \begin{pmatrix} -(\cos\phi + \sin\lambda)\sin\theta \\ \cos\epsilon(\sin\phi - \cos\lambda)\sin\theta - \omega t \sin\epsilon \cos\theta \\ \sin\epsilon(\sin\phi - \cos\lambda)\sin\theta + \omega t \cos\epsilon \cos\theta \end{pmatrix} . \tag{9}$$

At the end of the first pulse, the change in the spacecraft's velocity is  $(\phi = \pi + \lambda)$ :

$$\Delta \bar{v} = \frac{f}{m\omega} \begin{pmatrix} -(\sin\lambda - \cos\lambda) & \sin\theta \\ -(\sin\lambda + \cos\lambda) & \cos\epsilon\sin\theta - \frac{\pi}{2} & \sin\epsilon\cos\theta \\ -(\sin\lambda + \cos\lambda) & \sin\epsilon\sin\theta + \frac{\pi}{2}\cos\epsilon\cos\theta \end{pmatrix}.$$
 (10)

Integrating Equations (9) once again we obtain

$$\begin{pmatrix} \Delta X_{Q} \\ \Delta Y_{Q} \\ \Delta Z_{Q} \end{pmatrix} = -\frac{f}{m\omega^{2}} \begin{pmatrix} (\sin\phi - \cos\lambda + \omega t \sin\lambda) & \sin\theta \\ (\cos\phi + \sin\lambda + \omega t \cos\lambda) & \cos\epsilon & \sin\theta + \frac{(\omega t)^{2}}{2} & \sin\epsilon & \cos\theta \\ (\cos\phi + \sin\lambda + \omega t & \cos\lambda) & \sin\epsilon & \sin\theta - \frac{(\omega t)^{2}}{2} & \cos\epsilon & \cos\theta \end{pmatrix} . (11)$$

At the end of the first pulse the change in the spacecraft's position is

$$\Delta_{r}^{-} = -\frac{f}{m\omega^{2}} \begin{pmatrix} (-\sin\lambda - \cos\lambda + \frac{\pi}{2} \sin\lambda) & \sin\theta \\ (-\cos\lambda + \sin\lambda + \frac{\pi}{2} \cos\lambda) & \cos\epsilon & \sin\theta + \frac{\pi^{2}}{8} \sin\epsilon & \cos\theta \\ (-\cos\lambda + \sin\lambda + \frac{\pi}{2} \cos\lambda) & \sin\epsilon & \sin\theta - \frac{\pi^{2}}{8} \cos\epsilon & \cos\theta \end{pmatrix}.$$
 (12)

#### Geometry of Motion

In order to interpret the motion of the center of mass, described by Equations (11), we shall assume that, without any loss of generality,  $X_Q(0) = Y_Q(0) = Z_Q(0) = 0, \text{ and } \lambda = -\frac{\pi}{2}. \text{ Then } \phi = \omega t, \text{ and Equations (11) yield}$ 

$$\begin{pmatrix} X_{Q} \\ Y_{Q} \\ \end{pmatrix} = \frac{f}{m\omega^{2}} \begin{pmatrix} (\phi - \sin\phi) \sin\theta \\ \\ (1 - \cos\phi) \cos\epsilon \sin\theta - \frac{\phi^{2}}{2} \sin\epsilon \cos\theta \\ \\ (1 - \cos\phi) \sin\epsilon \sin\theta + \frac{\phi^{2}}{2} \cos\epsilon \cos\theta \end{pmatrix}$$

In the space-fixed, ecliptic coordinates this motion is expressed by

$$\begin{cases} X = k(\phi - \sin\phi), \\ Y = k(1 - \cos\phi), \\ Z = k_1 \phi^2, \end{cases}$$

where  $\phi = \omega t$ , and

$$k = \frac{f \sin \theta}{m\omega^2}, \quad k_1 = \frac{f \cos \theta}{2m\omega^2}.$$

It is obvious that the center of mass moves on a cycloidal cylinder (looking more like wavy roofing material), parallel to the Z-axis, and proportionally to  $t^2$  (or  $\phi^2$ ) up or down the surface of the cylinder, depending on whether f is positive or negative. The radius of the generating circle of the cycloid is

$$k = \frac{f \sin \theta}{2} .$$

The rate at which the center of mass moves in Z-direction (out of the ecliptic plane) is

$$\frac{f \cos \theta}{m} t$$
,

and it is obvious that for  $\theta$  = 90° the center of mass remains in the ecliptic plane. With  $\theta$  = 38°, m = 356.9 kg, and f = 1 Nt, this rate is given in km/sec by

$$2.2 \times 10^{-6} t$$

where t is in seconds of time. This means that, due to the action of only one pulse, the spacecraft will move 47.5 m out of the ecliptic plane in one day.

#### The Motion During One Cycle (16 Pulses)

It is obvious that, during one complete cycle of sixteen pulses, the disturbing effects will accumulate. For instance, if  $\Delta \bar{r}$  and  $\Delta \bar{v}$  [given respectively by Equations (12) and (10)] are respective changes in the position of the spacecraft and its velocity at the end of the first pulse lasting one quarter of a second, the total change in  $\bar{r}$  and  $\bar{v}$  after one second will be

$$\Delta \bar{r}_1 = \Delta \bar{r} + \frac{3}{4} \Delta \bar{v}$$

$$\Delta \bar{v}_1 = \Delta \bar{v}$$
.

At the end of the second firing (second cycle) the total changes in  $\bar{r}$  and  $\bar{v}$  will be

$$\Delta \bar{r}_1 + \Delta \bar{r} + \frac{1}{4} \Delta \bar{v} = 2\Delta \bar{r} + \Delta \bar{v}$$

$$\Delta \bar{v} + \Delta \bar{v} = 2\Delta \bar{v} ,$$

because the effects of the jet force during each pulse are the same if it is the same type of maneuver. After two seconds these changes will be

$$\Delta \overline{r}_{2} = 2\Delta \overline{r} + \Delta \overline{v} + \frac{3}{4} \Delta \overline{v} = 2\Delta \overline{r} + (1 + 2\frac{3}{4}) \Delta \overline{v}$$

$$\Delta \overline{v}_{2} = 2\Delta \overline{v}$$

etc. For instance,

$$\Delta \overline{r}_{3} = 3\Delta \overline{r} + (1 + 2 + 3.\frac{3}{4}) \Delta \overline{v}$$

$$\Delta \overline{v}_{3} = 3\Delta \overline{v} ,$$

$$\Delta \overline{r}_{4} = 4\Delta \overline{r} + (1 + 2 + 3.\frac{3}{4}) \Delta \overline{v} ,$$

$$\Delta \overline{v}_{4} = 4\Delta \overline{v} .$$

At the end of the n-th second the changes in  $\bar{r}$  and  $\bar{v}$  are

$$\Delta \bar{r}_{n} = n\Delta \bar{r} + \frac{n(2n+1)}{4} \Delta \bar{v}$$
 (13)

$$\bar{\Delta v}_{n} = n \Delta \bar{v} \tag{14}$$

where  $\Delta \bar{r}$  and  $\Delta \bar{v}$  are given by Equations (12 and (10), respectively. For n = 16 (end of one cycle)

$$\Delta \bar{r}_{c} = 16 \Delta \bar{r} + 132 \Delta \bar{v}$$

$$\Delta \bar{v}_{c} = 16 \Delta \bar{v} . \qquad (15)$$

#### One Example

To illustrate the solution and to give a physical interpretation of the effects of the maneuvers on the position and velocity of the spacecraft's center of mass, we shall calculate the changes  $\Delta \bar{r}_c$  and  $\Delta \bar{v}_c$  after one complete cycle of firings. To simplify the calculations we shall assume that  $\lambda$  = 0 and obtain  $\Delta \bar{r}$  and  $\Delta \bar{v}$  respectively from Equations (12) and (10) (for the obliquity of the ecliptic we shall take the value  $\epsilon$  = 23° 4458). We find

$$\Delta \bar{r} = -\frac{f}{m\omega^2} \qquad (\frac{\pi}{2} - 1) \cos \epsilon \sin \theta + \frac{\pi^2}{8} \sin \epsilon \cos \theta$$

$$(\frac{\pi}{2} - 1) \sin \epsilon \sin \theta - \frac{\pi^2}{8} \cos \epsilon \cos \theta$$

or

$$\Delta \bar{r} = \begin{pmatrix} 4.3695 \\ -5.0335 \\ 5.3378 \end{pmatrix} \times 10^{-8} \text{ km}$$

and

$$\Delta \bar{v} = \frac{f}{m\omega} \begin{pmatrix} \sin\theta \\ -\cos\epsilon \sin\theta - \frac{\pi}{2} \sin\epsilon \cos\theta \\ -\sin\epsilon \sin\theta + \frac{\pi}{2} \cos\epsilon \cos\theta \end{pmatrix}$$

or

$$\Delta \bar{v} = \begin{pmatrix} 2.7455 \\ -4.7150 \\ 3.9717 \end{pmatrix} \times 10^{-7} \text{ km/sec.}$$

Hence,

$$\Delta \bar{r}_{c} = \begin{pmatrix} 3.6940 \\ -6.3043 \end{pmatrix} \times 10^{-5} \text{ km}$$
5.3280

$$\Delta \bar{v}_{c} = \begin{pmatrix} 4.3928 \\ -7.5440 \\ 6.3547 \end{pmatrix} \times 10^{-6} \text{ km/sec}$$

and,

$$|\Delta \bar{\mathbf{r}}| = 9.043 \times 10^{-2} \text{ m}, \qquad |\Delta \bar{\mathbf{v}}| = 1.080 \times 10^{-2} \text{ m/sec.}$$

The final velocity at the end of only one cycle would produce a change in the spacecrait's position of 0.93 km in one day.

#### The Computer Program

A computer program, based on an unperturbed elliptical orbit of Helios, is written in such a manner that it yields the heliocentric ecliptic longitude of the spacecraft,  $\lambda$ , which appears in Equations (10) and (12). The program computes the accumulated effects of the disturbing force in the heliocentric position and velocity of the spacecraft during first four maneuvers and can be extended to as many maneuvers as desirable. The cumulative effects on Helios A during the four maneuvers listed in Table 1 are given in Table 3. The complete listing of the program is given in the Appendix.

# The Implementation of the Perturbative Acceleration in the Orbit Determination Program

The disturbing effects in the position and velocity of the spacecraft, given by Equations (15), cannot be implemented into the Orbit Determination Program, therefore, the accelerations, producing these effects, should be included directly in the input data block of the program. The step-function accelerations, given by the three equations following Equations (8) cannot be introduced into the program either, as the program does not provide a mathematical model for such accelerations.

There is, however, no constant acceleration,  $\bar{a}$ , which, at the end of one cycle of 16 seconds, would produce changes  $\Delta \bar{r}_c$  and  $\Delta \bar{v}_c$ , specified by Equations (15). If there was such an acceleration, then the equations

$$\Delta \bar{r}_c = \frac{1}{2} \bar{a} T^2$$
,

Table 3. Accumulated disturbing effects in coordinates during maneuvers

```
MANUEUVRE NO.
                  1
       DX =
                -27.36286995 FM
       DY =
                -22.43863860 KM
       04 =
                 34.05021238 KM
      DVA =
                  -. 0000121n K4/SEC
      DVY =
                  -. 000000961 KM/SEC
      DVZ =
                  .0000150a KM/5EC
MANDEUVRE NO.
       UX =
                 10.78901064 KM
       DY =
                  2.24149281 KM
       DZ =
                -17.35809875 KM
      DVX =
                  -. 00001119 KM/5EC
      DVY =
                  -. DODODOZOA KM/SEC
      DVZ =
                   .DDDD183/ KM/SEC
MANGEUVRE NO.
       DX =
                  -. 32380591 KM
       DY =
                 -3.15407485 KM
       DZ =
                -15.28499794 KM
      DVX =
                  -. 00000183 KM/SEC
      DVY =
                  -. 00001524 KM/SEC
      DVZ =
                  -.00007400 KM/SEC
MA .. TUVRE NO.
       DX =
                  9.69688976 KM
       DY =
                -12.93451202 KM
       0 Z =
                 12.74077559 KM
                    . U0001528 KM/SEC
      DVX =
                  -.00002034 KM/SEC
      DVY =
      DVZ =
                   .00002006 KM/SEC
```

and

$$\Delta \bar{v}_{c} = \bar{a}T$$

with  $\Delta \bar{r}_c$  and  $\Delta \bar{v}_c$  given, and T = 16 sec, should be satisfied simultaneously. This, of course, is impossible, except in the case when  $2\Delta \bar{r}_c = T\Delta \bar{v}_c$ , or  $\Delta \bar{r}_c = 8\Delta \bar{v}_c$ , which is out of the question.

Since the orbit determination program provides an option for occasional accelerations, the problem can be solved in the following manner. Let  $\bar{a}_1$  be the acceleration given to the center of mass of the spacecraft at the beginning of a certain cycle, and  $\bar{a}_2$  be the acceleration added in the middle of the cycle, i.e. eight seconds after the first acceleration,  $\bar{a}_1$ , has started acting on the body; both  $\bar{a}_1$  and  $\bar{a}_2$  should be constant vectors and their combined actions should produce exactly the changes in the position and velocity of the spacecraft,  $\Delta \bar{r}_1$  and  $\Delta \bar{v}_2$  respectively, at the end of the cycle.

Equations describing this dynamic event are

$$\frac{1}{2}\bar{a}_1T^2 + \frac{1}{2}\bar{a}_2(\frac{T}{2})^2 = \Delta \bar{r}_c$$
,

and

$$\bar{a}_1 T + \bar{a}_2 (\frac{T}{2}) = \Delta \bar{v}_c$$
,

where T = 16 sec., or

$$128 \bar{a}_{1} + 32 \bar{a}_{2} = \Delta \bar{r}_{c} ,$$

$$16 \bar{a}_{1} + 8 \bar{a}_{2} = \Delta \bar{v}_{c} .$$

JPL Technical Memorandum 33-785

Solving for  $\overline{a}_1$  and  $\overline{a}_2$  we find

$$\bar{a}_1 = \frac{1}{64} \left( \Delta \bar{r}_c - 4 \Delta \bar{v}_c \right) ,$$

$$\bar{a}_2 = -\frac{1}{32} \left( \Delta \bar{r}_c - 8 \Delta \bar{v}_c \right) .$$
(16)

Expressed in terms of changes  $\Delta \bar{r}$  and  $\Delta \bar{v}$ , respectively given by Equations (12) and (10),  $\bar{a}_1$  and  $\bar{a}_2$  become

$$\bar{a}_1 = 0.25 \, \Delta \bar{r} + 1.0625 \, \Delta \bar{v}$$
,  
 $\bar{a}_2 = -0.5 \, \Delta \bar{r} - 0.125 \, \Delta \bar{v}$ 

It is evident that  $\bar{a}_1$  and  $\bar{a}_2$  can be introduced arbitrarily at any time  $t_1$  and  $t_2$  such that

$$0 \leq t_1 \leq t_2 \leq T = 16 \text{ sec.}$$

In that case, the two starting equations would have the forms

$$\frac{1}{2} \bar{a}_{1} \tau_{1}^{2} + \frac{1}{2} \bar{a}_{2} \tau_{2}^{2} = \Delta \bar{r}_{c} ,$$

$$\bar{a}_{1} \tau_{1} + \bar{a}_{2} \tau_{2} = \Delta \bar{v}_{c} ,$$

where  $\tau_1 = T - t_1$ ,  $\tau_2 = T - t_2$ . From these equations we obtain  $(\tau_1 \neq \tau_2)$ 

$$\bar{a}_1 = \frac{2}{\tau_1 (\tau_1 - \tau_2)} (\Delta \bar{r}_c - \frac{\tau_2}{2} \Delta \bar{v}_c)$$
,

$$\bar{a}_2 = \frac{2}{\tau_2 (\tau_1 - \tau_2)} (\Delta \bar{r}_c - \frac{\tau_1}{2} \Delta \bar{v}_c) .$$

#### Accelerations In The Sun-Canopus Oriented System

The system of reference axes, used in the Orbit Determination Program, is generated by means of the spacecraft-Sun and the spacecraft-star Canopus

directions. This is, of course, a rotating, non-inertial system, since the spacecraft moves in its orbit around the Sun. The spacecraft-Canopus direction, for all practical purposes, can be considered space-fixed.

The three unit vectors along the axes of the system are given by (Reference 1)

$$\bar{e}_1 = \frac{(\bar{U}_c \times \bar{r}) \times \bar{r}}{r |\bar{U}_c \times \bar{r}|} ,$$

$$\bar{e}_2 = \frac{\bar{v}_c \times \bar{r}}{|\bar{v}_c \times \bar{r}|}$$
,

$$\bar{e}_3 = \frac{\bar{r}}{r}$$
,

where

$$\bar{U}_{C} = \begin{pmatrix} \cos\alpha & \cos\delta \\ \sin\alpha & \cos\delta \\ \sin\delta \end{pmatrix}$$

is the unit vector along the spacecraft-Canopus direction.  $\alpha$  and  $\delta$  are respectively the right ascension and declination of Canopus for 1950.0 and their respective values are

$$\alpha = 98.02255$$

$$\delta = -68.98877$$
.

The components of accelerations  $\bar{a}_i$ , i = 1, 2 in this system of reference are then given by

$$\bar{a}_{i} = \begin{pmatrix} \bar{a}_{i} \cdot \bar{e}_{1} \\ \bar{a}_{i} \cdot \bar{e}_{2} \\ \bar{a}_{i} \cdot \bar{e}_{3} \end{pmatrix} ;$$

their values for the first four maneuvers of the Helios A spacecraft, obtained from the previously mentioned computer program, are listed in Table 4.

Table 4. Helios attitude control maneuvers, disturbing accelerations

TIME	XACC	YACC	ZACC
52.31180525	.575255093-06	447974720-06	.115085824-06
52,31189775	985046418-07	.704783474-07	139059737-07
52.31805515	.575554921-06	443400040-06	.130297547-06
52.31814766	985416406-07	.779102773-07	105808558-07
52.50416660	575610756-06	.442227012-06	133985235-06
52,50425911	.985476305-07	777615581-07	.172306311-07
52,51041651	5/5612482-06	.442189908-06	134100214-06
52,51050901	.985478206-07	777568401-07	172509103-07
52,51666641	575614010-06	.442155951-06	134205596-06
52,51675892	.985479938-07	-,777525013-07	·172694923-07
52,52291632	575615545-06	.442122008-06	134310850-06
52,52300882	.985481474-07	777481901-07	·172880525-07
52,58055544	.575629727-06	-,441807515-06	•135281400-06 ••174591661-07
52,58064795	985496236-07	.777 <sup>U8</sup> U542-07	.135608921-06
52.59999990	.575634573-06	.776744020-07	175169310-07
52.60009241	985501449-07	- 441593770-06	.135936469-06
52.61944437 52.61953688	.575639426-06 985536441-07	776807152-07	175746884-07
52.63888884	.575644314-06	441486392-06	.136264097-06
52.63898134	985511663-07	776669475-07	176324702-07
32,030,0134		• • • • • • • • • • • • • • • • • • • •	
74.32986069	602996828-06	.277100138-07	424819373-06
74.32795319	.102998788-06	107984350-07	.730071035-07
74.33982563	603019330-06	273417946-07	424811287-06
74.33991814	.103002/34-06	107351568-07	.730108711-07
74.34886551	603039759-06	.270075677-07	424803687-06
74.34895802	.103006315-06	106/77183-07	•730142435-07 ••424795612-06
74,35813610	603000677-06	.264645799-07 106187721-07	.730176604-07
74.35822868	•103009971-06	100.07/21-07 250888839-07	.424755154-06
74.40068245	.603156799-06 103026809-06	.103479314-07	730328003-07
74.40077496 74.42134190	.603203482-06	243226767-07	. 424733425-06
74.42143440	103034999-n6	.102161923-07	730397938-07
74.59218693	.603584662-06	179575945-07	. 424501547-06
74.59227943	103102759-06	.912110032-08	730851339-07
74.63725662	.603691561-06	162700431-07	. 424424655-06
74.63734913	103120626-06	.883055407-08	730995575-07
74.65961742	.603742130-06	154314943-07	.424384048-06
74.65970993	103129521-06	.868613892-08	731043386-07
74.67769623	.603783029-06	147528176-07	.424350006-06
74.67778873	1031367u8-u6	.856 <sup>9</sup> 24265-08	731079934-07

REPRODUCIBILITY OF THE ORIGINAL PAGE IS POOR

Table 4. (contd)

TIME	XACC	YACC	ZACC
71.31180477 91.31189728 91.32638836 91.32648087 91.35416503 91.35425854 91.36814785 91.36814785 91.38194370 91.38203621 91.37583302 91.39592552 91.40972137	601555598-06 .103082892-06601534701-06 .103078673-06601495628-06 .103070738-06601476444-06 .103066859-06601457515-06 .103063010-06601438792-06 .103059187-06	271563319-06 .513U31675-07 270527792-06 .511381604-07 268545U48-06 .508219724-07 267548671-06 .508629672-07 266549097-06 .505033713-07 265545978-06 .503431249-07	.330497748-06 528894981-07 .331383823-06 530572310-07 .333063184-06 533754330-07 .333896658-06 535338915-07 .334731187-06 536918918-07 .335561047-06 538494787-07 .336387885-06
91.40981388  113.38140011 113.38149261 113.40277767 113.40287018 113.40972137 113.40981388 113.41675854 113.42361069 113.42370319 113.43055534 113.43064785 113.43750000 113.43759251 113.45138836	.103055419-06 .593715429-06 .593681576-06 .593680618-06 -102074916-06 .593669014-06 -102072748-06 .593657447-06 -102070597-06 .593645886-06 -102068436-06 .593634340-06 .593634340-06 .593622829-06 -593599822-06	.501823059-07  .203435000-06293520523-07 .204231828-06294939100-07 .204490313-06295399296-07 .204748542-06295859222-07 .205006710-06296318827-07 .205264602-06296778229-07 .205522344-06297237217-07206037210-06	540065939-07  .388535110-06 690883590-07  .388171621-06 690377711-07  .388053248-06 690212891-07  .387934794-06 69982400-07  .387697344-06 689716764-07  .387578428-06 689550985-07 387540208-06
113,45148087 113,45833302 113,45842552 113,46527767 113,46537018 113,47222137 113,47231388	.102059853-06 593588361-06 .102057718-06 593576914-06 .102055592-06 593565481-06 .102053459-06	.298154408-07 206294441-06 .298612615-07 206551398-06 .299070386-07 206808146-06 .299527718-07	.469218327-07 387220862-06 .689051687-07 387101416-06 .68884683-07 386981831-06 .688717412-07

#### APPENDIX

#### COMPUTER PROGRAM FOR CALCULATION OF HELIOS MANEUVERS

```
C THIS PROGRAMME COMPUTES PERTURBATIONS DUE TO ATTITUDE CONTROL MANOEUVRES.
   FOR ARBITRARY NUMBER OF FIRING CYCLES. EACH CYCLE CONSISTS OF 16 PULSES.
   LACH PULSE LASTING ONE QUARTER OF A SECOND. JET FORCE IS ONE NEWTON.
     C***
C
000
   NOMENCLATURE-
           Ax
                  = THE SEMI-MAJOR AXIS OF THE SPACECRAFT-S ELLIPTIC
0000
                    ORBIT (KM)
           ECC
                  = THE ORBITAL ECCENTRICITY OF THE SPACECRAFT
           EPSLN = OBLIGUITY OF THE ECLIPTIC
           INCL
                  = INCLINATION OF THE ORBITAL PLANE TO THE ECLIPTIC
                    PLANE OF 1950.0
000000000000000
           NODE
                  = NODAL ANGLE OF THE ORBITAL PLANE OF SPACECRAFT
           UMEGA
                 = ARGUMENT OF THE PERIAPSIS OF THE SPACECRAFT-S ORBIT
           M
                  = MEAN ANOMALY OF THE SPACECRAFT
                  = ECCENTRIC ANOMALY OF THE SPACECRAFT
           MISTART = MEAN ANOMALY OF THE SPACECRAFT AT THE TIME OF
                    INITIALIZATION T = TSTART
           TSTART = INITIAL TIME
                  = GRAVITATIONAL CONSTANT OF THE SUN
= MEAN ORBITAL MOTION OF THE SPACECRAFT
           GM.
           MEAN
           PER
                  = ORBITAL PERIOD OF THE SPACECRAFT
                  = NUMBER OF POINTS ON THE TRAJECTORY
           TSTEP = TIME STEP (DAYS)
C
      EPOCH = ISTART = 1974. DECEMBER 10. 00 HRS. 00 MIN. 00 SEC.
C
   SPECIFICATIONS-
      REAL MSTART, MZERO, MEAN, LONG , MASS, INCL, NODE
      DIMENSION X(40), Y(40), Z(40), R(40), TA(40), TIME(40), DAN(40), SAT(40),
     1 EM(40), SEC(40), LONG(40), MANUVR(4), DELX(40), DELY(40), DELZ(40),
     2 DELDX(40),DELDY(40),DELDZ(40),TAU(4),DX(4),DY(4),DZ(4),DDX(4),
     3 DDY(4),DDZ(4),XA(2,40),YA(2,40),ZA(2,40),XACC(80),YACC(80),
     4 ZACC(80), THALF(80)
C
      NAME' IST/INPUT/TSTART, TSTEP, DT, AX, ECC, OMEGA, NODE, INCL, AU, MZERO,
     1 N.TL CH.GM. DAY, PI, EPSLN. THETA, MASS, SAT, EM, SEC, ALPHA, DELTA
 7000 READ (5, INPUT)
      WRITE(6,1000)
      WRITE(6, INPUT)
      WRITE(6,3000)
 1000 FORMAT(1H1,1x,/)
 3000 FORMAT(1H1,4X////)
      RAD = 100./PI
 COMPUTATION OF UNPERTURBED POSITIONS OF THE SPACECRAFT
      EPSLN = EPSLN/RAD
      CE = COS(EPSLN)
      SE = SIN(EPSLN)
      ALPHA = ALPHA/RAD
```

```
DELTA = DELTA/RAD
     XCAN = CCS(ALPHA) *COS(DELTA)
     YCAN = SIN(ALPHA) *COS(DELTA)
     ZCAN = SIN(DELTA)
     MEAN = SGRT(GM/AX**3)
           = MEAN*DAY*RAD
     PER = 2.0*PI/(MEAN*DAY)
     MSTART = MZERO - DM*DT
     IF (MSTART.LT.O.) MSTART=MSTART+360.
2003 WRITE(6,2)
     WRITE(6,3)AX, ECC, OMEGA, MSTART, DM
                                          , PER . INCL , NODE
     WRITE(6,32)
   2 FORMAT(36x,-HELIOCENTRIC EQUATORIAL ORBITAL PARAMETERS OF THE -,
    * -SPACECRAFT-//)
   3 FORMAT(38X,-AX =-,E16.8,1X,-KM-/,37X,-ECC =-,F16.11/,35X,-QMEGA =-
       *F16.11.1X,-DEG-/,34X,-MSTART =-,F16.11.1X,-DEG-/.36X,-MEAN =-,
        F16.11,1X,-DEG/DAY-/,34X,-PERIOD =-,F16.11,1X,-DAYS-/,36X,
       -INCL =-,F16.11,1x,-DEG-/,36x,-NODE =-,F16.11,1x,-DEG-///)
  32 FORMAT(1H1,24X,-TIMES OF FIRING CYCLES-/)
     CALL VECTOR(INCL, NODE, OMEGA, PX, PY, PZ, QX, GY, QZ)
     MSTART = MSTART/RAD
     DO 1 I=1.N
     IF(I.LT.11)DAN(I)=52.
     IF(I.GE.11.AND.I.LT.21)DAN(I)=74.
     IF(I.GE.21.AND.I.LT.28)DAN(I)=91.
     IF(I.GE.28)DAN(I)=113.
     TIME(I) = ((SEC(I)/6C + EM(I))/60 + SAT(I))/24 + DAN(I)
     WRITE(6,5) TIME(I)
     IF(I.EQ.10.OR .I.EQ.20.OR .I.EQ.27) WRITE(6,6)
   5 FORMAT (25x,F15.10)
   6 FORMAT(1X,/)
   1 CONTINUE
     WRITE(6,33)PX,PY,PZ,QX,QY,QZ
  33 FORMAT(1H1////,12x,-P =-/,3(15x,F12.7/),//,12x,-Q =-/,3(15x,
    1 F12.7/),//)
     WRITE(6.7)
   7 FORMAT(1H1,/,3x,-EQUATORIAL COORDINATES OF HELIOS-///,4x,-TIME-,
    1 10X,-X(KM)-,12X,-Y(KM)-,12X,-Z(KM)-,12X,-R(KM)-,1CX,-TA-,9X,
    2 -LONG-/)
     DO 8 I=1.N
     CALL ORDIT(1, TIME(1), MSTART, AX, ECC, MEAN, PX, PY, PZ, QX, QY, QZ,
      x(I), y(I), Z(I), R(I), TA(I), LONG(I))
     WRITE(6,9) TIME(I), X(I), Y(I), Z(I), R(I), TA(I), LONG(I)
     IF(I.EG.10.CR .I.EG.20.CR .I.EG.27)WRITE(6,91)
   9 FORMAT (F9.2,4E17.7,2F12.3)
  91 FORMAT(2X./)
   8 CONTINUE
  CALCULATION OF MANOEUVRES
     THETA = THETA/RAD
     SPIN = 2.*PI
     AL = P1/2.
     BL = (PI**2)/8.
     CL = AL - 1.
     CT = COS(THETA)
     ST = SIN(THETA)
     WRITE(6,98)
  98 FORMAT (1H1,2X/)
     DO 11 I=1.N
     IF(I \cdot LT \cdot 11) MANUVR(1) = 1
```

```
IF(I \circ GE \circ 11 \circ AND \circ I \circ LT \circ 21) MANUVR(2) = 2
   IF(I.GE.21.AND.I.LT.28)MANUVR(3) = 3
   IF (I.GE. 28) MANUVR (4) - 4
   SIGN = 1.
   IF(I.GE.3.AND.I.LE.6)SIGN=-1.C
   IF(I.GE.11.AND.I.LE.14)SIGN=-1.0
   IF(I.GE.21.AND.I.LE.27)SIGN=-1.0
   IF(1.GE.35)SIGN=-1.0
   FORCE = SIGN+1.E-03
   WRITE(6,99) FORCE
99 FORMAT(10x,-FORCE =-,F9.3,1x,-KG.KM/SEC**2-)
   RCOEF = -FORCE/(MASS*SPIN**2)
   VCOEF = -SPIN*RCOEF
   LONG(I) = LONG(I)/RAD
   CX = CL*SIN(LONG(I)) - CCS(LONG(I))
   CYZ = CL*COS(LONG(I)) + SIN(LONG(I))
   CVX = -SIN(LONG(I)) + COS(LONG(I))
   CVYZ = -(SIN(LONG(I)) + COS(LONG(I)))
   DRX = RCOEF*CX*ST
   DRY = RCGEF*(CYZ*CE*ST + BL*SE*CT)
   DRZ = RCOEF*(CYZ*SE*ST - BL*CE*CT)
   DVX = VCOEF*CVX*ST
   DYY = VCOEF*(CVYZ*CE*ST - AL*SE*CT)
   DVZ = VCOEF*(CVYZ*SE*ST + AL*CE*CT)
   DELX(I) = 16.*DRX + 132.*DVX
   DELY(I) = 16.*DRY + 132.*DVY
   DELZ(I) = 16.*DRZ + 132.*DVZ
   DELDX(I) = 16.*DVX
   DELDY(I) = 16.*DVY
   DELDZ(I) = 16.*DVZ
11 CONTINUE
   WRITE(6,15)
15 FORMAT(1H1,2x,-TIME-,9x,-Dx(KM)-,11x,-DY(KM)-,11x,-DZ(KM)-,11x,
  1 -DVX(KM/5)-,8X,-DVY(KM/5)-,8X,-DVZ(KM/5)-/)
   DO 22 I=1.N
   wRITE(6,16)TIME(1),DELX(1),DELY(1),DELZ(1),DELDX(1).DELDY(1).
     DELDZ(I)
   IF(I.EQ.10.OR .I.EQ.20.OR .I.EQ.27)WRITE(6,20)
16 FORMAT (1x, F7.2, 6F17.9)
20 FORMAT (1X./)
22 CONTINUE
    WRITE(6,4)
 4 FORMAT(1H1,//,10x,-ACCUMULATED DISTURBING EFFECTS IN COORDINATES -
  1 -DURING MANOEUVRES-//)
   00 10 J=1,4
   Dx(J) = 0.0
   DY(J) = 0.0
   DZ(J) = 0.0
   DDX(J) = 0.0
   DDY(J) = 0.0
   DDZ(J) = 0.0
   IF(J.EG.1) ISTRT=1
   IF (J.EG.2) ISTRT=11
   IF (J.EG.3) ISTRT=21
   IF (J.EQ.4) ISTRT=28
   IF(J.EQ.1)NFIN=10
   IF (J.EQ.2) NF IN= 20
   IF (J.EG.3) NF IN=27
```

```
IF (J. EG. 4) NF IN= 38
   TAU(J) = TIME(NFIN)
   DO 12 I=ISTRT,NFIN
   K = I - ISTRT + 1
   L = I-1
   DDX(J) = DDX(I) + DELDX(I)
   DDY(J) = DDY(J, + DELDY(I)
   DDZ(J) = DDZ(J) + DELDZ(I)
   IF(K.EQ.1)GO TO 14
   DX(J) = DX(J) + DELX(I)
   DY(J) = DY(J) + DELY(I)
   DZ(J) = DZ(J) + DELZ(I)
   GO TO 12
14 DELT = (TAU(J) - TIME(L))*DAY
   DX(J) = DX(J) + DELX(I) + DELDX(I)*DELT
   DY(J) = DY(J) + DELY(I) + DELDY(I)*DELT
   DZ(J) = DZ(J) + DELZ(I) + DELDZ(I)*DELT
12 CONTINUE
   wRITE(6,17)MANUVR(J),DX(J),DY(J),DZ(J),DDX(J),DDY(J),DDZ(J)
17 FORMAT(7x,-MANOEUVRE NO.-,15//,14x,-Dx =-,F16.8,1x,-KM-/,14x,
  1 - DY = -,F16.8,1x,-KM-/,14x,-DZ = -,F16.8,1x,-KM-/,13x,-DVX = -,
  2 F16.8,1x,-KM/SEC-/,13x,-DVY =-,F16.8,1x,-KM/SEC-/,13x,-DVZ =-,
  3 F16.8.1X.-KM/SEC-///)
10 CONTINUE
CALCULATION OF ACCELERATIONS WHICH SHOULD BE IMPLEMENTED INTO THE ORBIT
DETERMINATION PROGRAMME
   20 18 I=1.N
           = (DELX(I) - 4.*DELDX(I))/64.
   Alx
   AlY
           = (DELY(I) - 4.*DELDY(I))/64.
   AlZ
           = (DELZ(I) - 4.*DELDZ(I))/64.
           = (-DELX(I) + 8.*DELDX(I))/32.
   AZX
   A2Y
           = (-DELY(1) + 8.*DELDY(1)//32.
           = (-DELZ(1) + 8.*DELDZ(1))/32.
   AZZ
TRANSFORMATION INTO THE XSTAR, YSTAR, USP SYSTEM
   JDR = XCAN*X(I) + YCAN*Y(I) + ZCAN*Z(I)
   AMAG = SQRT(R(I)**2 - UDR**2)
   UXX = YCAN*Z(I) - ZCAN*Y(I)
   UXY = ZCAN*X(I) - XCAN*Z(I)
   UXZ = XCAN*Y(I) - YCAN*X(I)
   E1X = (X(I)*JDR - XCAN*R(I)**2)/(AMAG*R(I))
   E1Y = (Y(I)*UDR - YCAN*R(I)**2)/(AMAG*R(I))
   E1Z = (Z(I)*UDR - ZCAN*R(I)**2)/(AMAG*R(I))
   EZX = UXX/AMAG
   EZY = LXY/AMAS
   E2Z = UXZ/AMAG
   E3X = X(I)/R(I)
   E3Y = Y(I)/R(I)
   E3Z = Z(I)/R(I)
   XA(1,I) = A1X*E1X + A1Y*E1Y + A1Z*E1Z
   YA(1,I) = A1X*E2X + A1Y*E2Y + A1Z*E2Z
   ZA(1,1) = A1X*E3X + A1Y*E3Y + A1Z*E3Z
   XA(2,I) = A2X*E1X + A2Y*E1Y + A2Z*E1Z
   YA(2*I) = A2X*E2X + A2Y*E2Y + A2Z*E2Z
   ZA(2 \cdot 1) = A2X*E3X + A2Y*E3Y + A2Z*E3Z
   I1 = 2 \times I - 1
   12 = 2*1
   XACC(11) = XA(1.1)
   XACC(I2) = XA(2.I)
   YACC(II) = YA(I \cdot I)
```

```
YACC(12) = YA(2,1)
     ZACC(II) = ZA(I,I)
     ZACC(12) = ZA(2,1)
     THALF(I1) = TIME(I)
     THALF(12) = TIME(1) + 8./DAY
  18 CONTINUE
     NT2 = 2*N
     DO 21 I=1,NT2
     IF(I.EQ.1.OR.1.EQ.41) WRITE(6,19)
     WRITE(6,23)THALF(1), XACC(1), YACC(1), ZACC(1)
     IF(1.EQ.20.OR .I.EQ.40.OR .I.EQ.54)WRITE(6,24)
  19 FORMAT(1H1,//,3X,-DISTURBING ACCELERATIONS IN THE XSTAR, YSTAR, -
    1 -USP SYSTEM IN KM/SEC++2-//,3X,-TIME-,15X,-XACC-,13X,-YACC-,13X,
    2 -ZACC-/)
  23 FORMAT(1x,F13.8,3E17.9'
  24 FORMAT(1X./)
  21 CONTINUE
     WRITE (6,2000)
2000 FORMAT(1H1,2X/)
     STOP
     END
```

```
SUBROUTINE VECTOR (INCL, NODE, OMEGA, PX, PY, PZ,QX,G,,QZ)
THIS SUBROUTINE COMPUTES THE COMPONENTS OF VECTORS P. Q. AND R.
    REAL INCL.NODE
DATA PI/3.141592654/
    RAD = 180./PI
    INCL = INCL/RAD
    NODE - NODE/RAD
    CMEGA - OMEGA/RAD
    CI = COS(INCL)
    SI = SIN(INCL)
    CN = COS(NODE)
    SN = SIN(NODE)
    CO = COS(C GA)
    SO = SINTOMEGA)
    PX = CN+CO - SN+SO+CI
    PY = SN*CO + CN*SO*CI
    PZ = 50*51
    QX = -CN+SO - SN+CO+C!
    QY = -SN*SO + CN*CO*CI
    QZ = CO*SI
    RX = SN*SI
RY = -CN*SI
    RZ = CI
    RETURN
    END
```

```
SUBROUTINE ORBIT (BODY, TIME, MSTART, SMA, ECC, MEAN, PX, PY, PZ, QX,
     1 QY,QZ,X,Y,Z,R,T,LONG)
C THIS SUBROUTINE COMPUTES HELIOCENTRIC EQUATORIAL COORDINATES OF HELIOS
      INTEGER BODY
      REAL M.MSTART, MEAN, LONG
      DATA TU/5.0/,PI/3.141592654/,DAY/86400./,EPSLN/23.445789/
      RAD = 180./PI
      EPSLN = EPSLN/RAD
      ETA = SGRT((1.0+ECC)/(1.0-ECC))
      M = MSTART + (TIME-TO)*DAY*MEAN
      E = ANOM(ECC.M)
      R = SMA*(1.-ECC*COS(E))
      T = 2.*ATAN(ETA*TAN(E/2.))
      IF(T.GE.2.*PI)T=T-2.*PI
      XORB = R*COS(T)
      YORB = R*SIN(T)
      T = T*RAD
      IF(T.LT.0.) T=T+360.
      X = PX*XORB + QX*YORB
      Y = PY*XCRB + GY*YORB
      Z = PZ*XORB + GZ*YORB
      SE = SIN(EPSLN)
      CE = COS(EPSLN)
C ECLIPTIC LONGITUDE OF HELIOS
      CL = X/R
      SL = (Y*CE + Z*SE)/R
      LONG = ATAN2(SL,CL)
      LONG = LONG*RAD
      IF (LONG.LT.O.)LONG=LONG+360.
      RETURN
      END
```

FUNCTION ANOM(ECC.M)

C THIS FUNCTION SUBROUTINE SOLVES THE KEPLER-S EQUATION BY ITERATIONS

REAL M

DATA EPS/.000005/

ANOM = M

2 ANOM = M + ECC\*SIN(ANOM)

TEST = ANOM - M - ECC\*SIN(ANOM)

IF(ABS(TEST).GT.EPS)GO TO 2

1F(ABS(TEST) . LE . EPS)RETURN

END

```
SINPUT
 ECC = .521807390542
 TSTART = U.U
 TSTEP = 86400.0
 AX = .96001973563E+08
 ALPHA = 98.02255
 DELTA = -68.98877
 EPSLN = 23.44578889
 JMEGA = 257.444938784
 INCL = 23.4469891399
 NOUE = .0613804058997
 MZERO = 71.6320412103
 TLNCH = -5.
 AJ = .1495978930E+9
GM = .132712499390602500E+12
 MASS = 356.9
 DT = 128.2152777778
 DAY = 86400.
 in = 30
 PI = 3.141592654
 THETA = 30.
 SAT = 2*7.,4*12.,13.,2*14.,15.,7.,3*8.,9.,10.,14.,2*15.,16.,2*7.,2*8.,
   6*9.,5*10.,3*11.
 EM = 29.938.96.915.924.933.956.924.952.920.955.99.922.935.936.96.912.9
   17.,49.,15.,29.,50.,30.,50.,10.,30.,50.,9.,40.,50.,0.,10.,20.,30.,
   50.,0.,10.,20.
 SEC = 11*0.,21.,22.,43.,59.,44.,45.,39.,51.,53.,7*0.,13.,10*0.
 SEND
-FIN
                                                                                NIF-
```

JPL Techanical Memorandum 33-785

#### References

1. Georgevic, R. M., Mction of the Sun-Canopus Oriented Attitude Control Reference Frame of the Mariner Venus/Mercury Spacecraft, Technical Memorandum 391-429, Jet Propulsion Laboratory, Pasadena, Calif., March 30, 1973 (an internal document).

#### Bibliography

1. Georgevic, R. M., Rigid Body Dynamics, Lecture N. tes, 1967-1968.